# Multiple Solutions and Regularity of H-systems 

Libin Mou and Paul Yang

April 1994


#### Abstract

The main result of this paper proves the existence of multiple solutions to a class of generalized constant mean curvature equations, called H -systems. Also contained is a regularity for conformal nharmonic maps.


## 1 Introduction

In this paper, we consider some systems of the form

$$
\begin{equation*}
\operatorname{div}\left(|\nabla u|^{n-2} \nabla u\right)=f(u, \nabla u), \tag{1}
\end{equation*}
$$

where $u \in W^{1, n}\left(\Omega, R^{k}\right), n, k \geq 2 ; \Omega \subset R^{n}$ is a bounded smooth domain, and $f: R^{k} \times R^{n k} \rightarrow R^{k}$ is a smooth function. We assume

$$
\begin{equation*}
|f(u, \nabla u)| \leq \Lambda|\nabla u|^{n}, \tag{2}
\end{equation*}
$$

for some constant $\Lambda>0$ that may depend on $u$.
A well-known example of (1) is the $n$-harmonic map equation. Let $(N, h) \hookrightarrow$ $R^{k}$ be a $C^{\infty}$ compact Riemannian submanifold. An $n$-harmonic map $u: \Omega \rightarrow$ $N$ is a critical point of the $n$-energy $\int_{\Omega}|\nabla u|^{n} d x$ in the space of functions $u \in W^{1, n}\left(\Omega, R^{k}\right)$ with $u(x) \in N$ for a.e. $x \in \Omega$. The equation for n-harmonic maps is

$$
\begin{equation*}
\operatorname{div}\left(|\nabla u|^{n-2} \nabla u\right)=|\nabla u|^{n-2} Q(u, \nabla u), \tag{3}
\end{equation*}
$$

where $Q(u, \cdot)$ is the trace of the second fundamental form of $N$ at $u(x) \in N$; $Q(u, \nabla u)$ is quadratic in $\nabla u$.

There is a vast literature on the regularity and partial regularity of solutions to harmonic (or $p$-harmonic) map type equations; see [4][11][13][15][17][19][20][24][26][30] and other references therein.

Our interest in this paper is mainly on the H -systems in higher dimensions. Suppose $u \in W^{1, n}\left(\Omega, R^{n+1}\right), u=\left(u^{1}, \ldots, . u^{n+1}\right)$. Then the cone generated by the image $u(\Omega)$, with vertex being the origin of $R^{n+1}$, has a welldefined volume

$$
V(u)=\frac{1}{n+1} \int_{\Omega} u \cdot u_{1} \wedge \cdots \wedge u_{n}
$$

see [24]. Here $u_{1} \wedge \cdots \wedge u_{n}$ is the cross product of the partial derivatives $u_{1}, \ldots, u_{n}$, which can be described as follows. For any vector $v \in R^{n+1}$,

$$
v \cdot u_{1} \wedge \ldots \wedge u_{n}=\left|\begin{array}{cccc}
v^{1} & v^{2} & \ldots & v^{n+1} \\
u_{1}^{1} & u_{1}^{2} & \ldots & u_{1}^{n+1} \\
\ldots & \ldots & \ldots & \ldots \\
u_{n}^{1} & u_{n}^{2} & \ldots & u_{n}^{n+1}
\end{array}\right| .
$$

Consider the minimization problem

$$
\begin{equation*}
\min \int_{\Omega}|\nabla u|^{n}, u=\eta \text { on } \partial \Omega, V(u)=c, \tag{4}
\end{equation*}
$$

for a given $\eta \in W^{1, n}\left(\Omega, R^{n+1}\right)$ and a constant $c$. A critical point of (4) is called an $n$-harmonic map with prescribed volume; it satisfies

$$
\begin{equation*}
\operatorname{div}\left(|\nabla u|^{n-2} \nabla u\right)=H u_{1} \wedge \cdots \wedge u_{n}, u=\eta \text { on } \partial \Omega, \tag{5}
\end{equation*}
$$

where $H$ is the Lagrange multiplier.
When $n=2$, (5) becomes

$$
\begin{equation*}
\Delta u=H u_{1} \wedge u_{2} \tag{6}
\end{equation*}
$$

A conformal solution of (6) represents a surface of constant mean curvature; see, e.g., [28] [31]. The existence of solutions and multiple solutions of (6) were established in many works, including [6] [21] [27] [29] [31] [33]. In Theorems 5 and 12 below, we prove for relatively small $H$ and boundary data, there is a solution of least energy-the small solution, and there is a
large solution, with the same boundary data. This generalizes the early work of Hildebrandt [21], Brezis and Coron [6] and Struwe [29] for $n=2$.

For the regularity of (2-)harmonic maps $u$ on a domain $\Omega \subset R^{2}$ (or a smooth surface), Heléin [19] proved their $C^{\infty}$ regularity. Assuming $u$ is conformal, or stationary or energy minimizing, Morrey[23], Grüter[16] and Schoen[25] established the regularity of $u$ earlier. For the $H$-system (6) with constant $H$, Wente [31] showed that any solution of (6) is analytic. Grüter [16] proved the $C^{1, \alpha}$ regularity ( $0<\alpha<1$ ) of conformal solutions to (6), where $H$ may depend on $u$; same result was obtained later by Bethuel [5] assuming that $|D H(u)|$ is bounded. Wente's result was generalized to (5) in $[10][24]$, which implies that all solutions of (5) are $C^{1, \alpha}$ regular. In this paper, we prove the $C^{1, \alpha}$ regularity of conformal solutions to (1), which generalizes the work of Grüter [16]. In particular, conformal n-harmonic maps from $\Omega \subset R^{n}$ (or an n-manifold) are $C^{1, \alpha}$, and conformal solutions of (5) with bounded $H=H(u)$ are also $C^{1, \alpha}$. Unlike in two dimension, one cannot reparametrize a solution to obtain conformality; so the conformality condition for solutions to (1) is fairly strong. It is conjectured that all nharmonic maps and solutions to (5) with bounded $H=H(u)$ are $C^{1, \alpha}$. Generally speaking, $C^{1, \alpha}$ regularity is optimal for solutions of (1)as shown by examples in [22].

## 2 Existence of Solutions to H-systems

For any $u \in W^{1, n}\left(\Omega, R^{n+1}\right)$, the image $u(\Omega)$ is a generalized "hypersurface" with area

$$
A(u)=\int_{\Omega} J(u) d x, J(u)=\left|u_{1} \wedge \cdots \wedge u_{n}\right|,
$$

where $J(u)$ is the Jacobian of $u$. Note that

$$
\begin{align*}
\left|v \cdot u_{1} \wedge \cdots \wedge u_{n}\right| & \leq|v|\left|u_{1}\right| \cdots\left|u_{n}\right| \\
& \leq|v|\left(\frac{\left|u_{1}\right|^{2}+\cdots+\left|u_{n}\right|^{2}}{n}\right)^{n / 2}  \tag{7}\\
& =|v| \frac{|\nabla u|^{n}}{\sqrt{n^{n}}},
\end{align*}
$$

and the equalities hold if and only if $u$ is conformal. Here we say that a function $u \in W^{1, n}\left(\Omega, R^{k}\right)$ is conformal if for some function $\lambda(x)$ and all $i, j=1, \ldots, n$,

$$
\begin{equation*}
u_{i} \cdot u_{j}=\lambda(x) \delta_{i j} . \tag{8}
\end{equation*}
$$

It follows from (7)

$$
\begin{equation*}
\left|u_{1} \wedge \cdots \wedge u_{n}\right| \leq \frac{|\nabla u|^{n}}{\sqrt{n^{n}}} \text { and } A(u) \leq \frac{1}{\sqrt{n^{n}}} \int_{\Omega}|\nabla u|^{n} \tag{9}
\end{equation*}
$$

and each of the equalities holds iff $u$ is conformal.
We now discuss some properties of the volume functional $V$.
First note that if $u=\left(u^{1}, \ldots, u^{n+1}\right)$ and $u^{1}=0$ on $\partial \Omega$, then for all $i=$ $1, \ldots, n$,

$$
\begin{equation*}
\int_{\Omega} u^{1} \frac{\partial\left(u^{2}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)}=(-1)^{i-1} \int_{\Omega} u^{i} \frac{\partial\left(u^{1}, \ldots, \hat{u^{i}}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)} \tag{10}
\end{equation*}
$$

and the volume $V$ can be written as

$$
\begin{equation*}
V(u)=\int_{\Omega} u^{1} \frac{\partial\left(u^{2}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)} . \tag{11}
\end{equation*}
$$

In fact, (10) follows by expanding the determinant $\frac{\partial\left(u^{2}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)}$ in the $i$-th column, using integration by parts together with the fact

$$
\sum_{\alpha=1}^{n} \frac{\partial}{\partial x^{\alpha}} \frac{\partial\left(u^{2}, \ldots, \hat{u}^{i}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, \hat{x^{\alpha}}, \ldots, x^{n}\right)}=0
$$

Expanding $V(u)$ in terms of $u^{1}, \ldots, u^{n+1}$ we get (11) by using (10).
As a consequence of (10) and isoperimetric inequality, we have
Proposition 1 If $u=\left(u^{1}, \ldots, u^{n+1}\right) \in W^{1, n}\left(\Omega, R^{1+n}\right)$ and $u^{1}=0$ on $\partial \Omega$, then for some constant $C_{1}$,

$$
\begin{equation*}
\left|\int u^{1} \frac{\partial\left(u^{2}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)}\right| \leq C_{1}\left\|\nabla u^{1}\right\|_{L^{n}(\Omega)} \cdots\left\|\nabla u^{n+1}\right\|_{L(\Omega)} \tag{12}
\end{equation*}
$$

Proof: We may assume that none of $u^{i}$ is constant (otherwise, the inequality is trivial), and that $\left\|\nabla u^{i}\right\|_{L^{n}(\Omega)}=1$ for all $i$ (by the homogeneity of (12) in $u^{i}$ ). Then (9) implies

$$
A(u) \leq \frac{1}{\sqrt{n^{n}}} \int_{\Omega}|\nabla u|^{n}=\left(\frac{n+1}{n}\right)^{n / 2} .
$$

Denote $v=\left(0, u^{2}, \ldots, u^{n+1}\right)$. Then $A(v) \leq A(u)$ and $V(v)=0$. So

$$
\left|\int_{\Omega} u^{1} \frac{\partial\left(u^{2}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)}\right|=|V(u)-V(v)|
$$

is the volume enclosed by the graphs of $u$ and $v$, whose area is $A(u)+A(v)$. By isoperimetric inequality (see [2], for example),

$$
|V(u)-V(v)| \leq \frac{1}{C}[A(u)+A(v)]^{(n+1) / n}
$$

where $C=(n+1) \omega_{n}^{\frac{1}{n}}$ and $\omega_{n}$ is the area of the unit $n$-sphere $S^{n}$. Therefore, for an absolute constant $C_{1}$,

$$
\left|\int_{\Omega} u^{1} \frac{\partial\left(u^{2}, \ldots, u^{n+1}\right)}{\partial\left(x^{1}, \ldots, x^{n}\right)}\right| \leq C_{1}
$$

which shows (12).
(12) implies the following corollaries.

Corollary 2 If $u \in W^{1, n}\left(\Omega, R^{n+1}\right)$ and $u^{i} \mid \partial \Omega=0$ for some $i=1, \ldots, n+$ 1 , then the functional $V$ is continuous at $u$ in the norm of $W^{1, n}\left(\Omega, R^{n+1}\right)$.

Corollary 3 Suppose $u, v \in W^{1, n}\left(\Omega, R^{n+1}\right)$ and $v=0$ or $u=0$ on $\partial \Omega$, then for some constant $C$,

$$
\begin{equation*}
\left|\int_{\Omega} v \cdot u_{1} \wedge \cdots \wedge u_{n}\right| \leq C\|\nabla v\|_{L^{n}(\Omega)}\|\nabla u\|_{L(\Omega)}^{n} \tag{13}
\end{equation*}
$$

Proof: Expand $\left|\int_{\Omega} v \cdot u_{1} \wedge \cdots \wedge u_{n}\right|$ in terms of $v^{1}, \ldots, v^{n+1}$ and apply (12) to each term.

We now derive a useful property of

$$
R(v, u)=\int_{\Omega} v \cdot u_{1} \wedge \ldots \wedge u_{n}
$$

Suppose $u, v, w \in W^{1, n}\left(\Omega, R^{n+1}\right), w=0$ or $v=0$ on $\partial \Omega$, and $u_{t}=u+t w$ for $0 \leq t \leq 1$. For a moment, suppose that $u, v, w \in C^{2}$. Then

$$
\begin{align*}
R(v, u+w)-R(v, u) & =\left.\int_{\Omega} v \cdot\left(u_{t}\right)_{1} \wedge \ldots \wedge\left(u_{t}\right)_{n}\right|_{0} ^{1} \\
& =\int_{\Omega} \sum_{i=0}^{n} v \cdot\left(u_{t}\right)_{1} \wedge \ldots \wedge w_{i} \wedge \ldots \wedge\left(u_{t}\right)_{n} \\
& =-\int_{\Omega} \int_{0}^{1} \sum_{i=0}^{n} w_{i} \cdot\left(u_{t}\right)_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge\left(u_{t}\right)_{n} \\
& =\int_{\Omega} \int_{0}^{1} \sum_{i=0}^{n} w \cdot\left(u_{t}\right)_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge\left(u_{t}\right)_{n} \\
& \sum_{j \neq i} \sum_{i=0}^{n} w \cdot\left(u_{t}\right)_{1} \wedge \ldots \wedge\left(u_{t}\right)_{j i} \ldots \wedge v_{i} \wedge \ldots \wedge\left(u_{t}\right)_{n} \\
& =\int_{\Omega} \int_{0}^{1} \sum_{i=0}^{n} w \cdot\left(u_{t}\right)_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge\left(u_{t}\right)_{n} . \tag{14}
\end{align*}
$$

Here we used the skew-symmetry of the cross product, which implies the term $\sum_{j \neq i} \sum_{i=0}^{n} \cdots=0$. It follows

$$
\begin{equation*}
|R(v, u+w)-R(v, u)| \leq C\|w\|_{\infty}\|\nabla v\|_{L^{n}}\left|\left\|\nabla u \left|+|\nabla w| \|_{L^{n}}^{n-1} ;\right.\right.\right. \tag{15}
\end{equation*}
$$

or

$$
\begin{equation*}
|R(v, u+w)-R(v, u)| \leq C\|\nabla v\|_{\infty}\|w\|_{L^{n}}\left|\|\nabla u|+| \nabla w\|_{L^{n}}^{n-1} .\right. \tag{16}
\end{equation*}
$$

The estimates (15) and (16) show that, in addition to the condition that $u, v, w \in W^{1, p}\left(\Omega, R^{n+1}\right)$, it is enough to assume $w \in C^{0}$ for (15) to hold, and $v \in W^{1, \infty}\left(\Omega, R^{n+1}\right)$ for (16).

Applying (14) to $u=0$ and $v, w \in W^{1, n}\left(\Omega, R^{n+1}\right)$ with $v$ or $w=0$ on $\partial \Omega$, we obtain

$$
\begin{align*}
\int_{\Omega} v \cdot w_{1} \wedge \cdots \wedge & w_{n}=\int_{\Omega} \int_{0}^{1} t^{n-1} d t \sum_{i=0}^{n} w \cdot w_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge w_{n}  \tag{17}\\
& =-\frac{1}{n} \sum_{i=1}^{n} \int_{\Omega} v_{i} \cdot w_{1} \wedge \cdots \wedge w_{i} \wedge \cdots \wedge w_{n} .
\end{align*}
$$

The equation (5) can be derived by using (17). We only need to calculate $\frac{d}{d t} V(u+t \phi)$ for any $\phi \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$. By (17)

$$
\begin{aligned}
& \frac{d}{d t} V(u+t \phi) \\
= & \frac{1}{n+1} \int_{\Omega} \phi \cdot u_{1} \wedge \cdots \wedge u_{n}+\frac{1}{n+1} \sum_{i=1}^{n} \int_{\Omega} u \cdot u_{1} \wedge \cdots \wedge \phi_{i} \wedge \cdots \wedge u_{n} \\
= & \frac{1}{n+1} \int_{\Omega} \phi \cdot u_{1} \wedge \cdots \wedge u_{n}-\frac{1}{n+1} \sum_{i=1}^{n} \int_{\Omega} \phi_{i} \cdot u_{1} \wedge \cdots \wedge u_{i} \wedge \cdots \wedge u_{n} \\
= & \int_{\Omega} \phi \cdot u_{1} \wedge \cdots \wedge u_{n} .
\end{aligned}
$$

The following is another property of $R$ that we prove by (17).
Theorem 4 Suppose that, as $m \rightarrow \infty, u^{m} \rightharpoonup u$ in $W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$, and either $v^{m} \rightarrow v$ in $W^{1, n}\left(\Omega, R^{n+1}\right)$ or $\left\|v^{m}-v\right\|_{\infty} \rightarrow 0$ with $v$ being continuous, then

$$
R\left(v^{m}, u^{m}\right) \equiv \int v^{m} \cdot u_{1}^{m} \wedge \ldots \wedge u_{n}^{m} \rightarrow R(v, u), \text { as } m \rightarrow \infty
$$

Proof: By (13) and the assumptions, we have

$$
\left|R\left(v^{m}, u^{m}\right)-R\left(v, u^{m}\right)\right| \leq\left\{\begin{array}{c}
\left\|\nabla v^{m}-\nabla v\right\|_{L^{n}(\Omega)} \\
\text { or }\left\|v^{m}-v\right\|_{\infty}
\end{array}\right\}\left\|\nabla u^{m}\right\|_{L^{n}(\Omega)}^{n} \rightarrow 0
$$

as $m \rightarrow \infty$. This implies that we may assume $v^{m} \equiv v$. Furthermore, we may assume that $v$ is $C^{2}$ by approximating $v$ by smooth functions in the norm of $W^{1, n}$, and in the norm of $C^{0}$ in case $v$ is continuous.

Now, because $u^{m} \rightharpoonup u$ in $W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$, we have $u^{m} \rightarrow u$ in $L^{n}$. By (17),

$$
\begin{aligned}
R\left(v, u^{m}\right) & =-\frac{1}{n} \sum_{i=1}^{n} \int_{\Omega} v_{i} \cdot\left(u^{m}\right)_{1} \wedge \cdots \wedge u_{i}^{m} \wedge \cdots \wedge\left(u^{m}\right)_{n} \\
& \rightarrow-\frac{1}{n} \sum_{i=1}^{n} \int_{\Omega} v_{i} \cdot u_{1} \wedge \cdots \wedge u_{i} \wedge \cdots \wedge u_{n}, \text { as } m \rightarrow \infty \\
& =\int_{\Omega} v \cdot u_{1} \wedge \cdots \wedge u_{i} \wedge \cdots \wedge u_{n}=R(v, u) .
\end{aligned}
$$

We now prove the existence of the small solutions.

Theorem 5 Suppose $\eta \in W^{1, n}\left(\Omega, R^{n+1}\right)$ and $0 \neq H$ is a constant satisfying

$$
\begin{equation*}
\|\eta\|_{\infty}|H| \leq \sqrt{n^{n}} \tag{18}
\end{equation*}
$$

Then the problem (5) has a solution $u$ that satisfies $\|u\|_{\infty} \leq\|\eta\|_{\infty}$.
Remark 6 The case $n=2$ of this theorem is due to Hildebrandt [21]; see also [20][27][31][32][10]. In the next section, we will show that if $\eta$ and $H$ are small enough, then the problem (5) has another "big" solution.

Remark 7 In general, a bound condition for $H$ like (18) is needed for the existence of a solution. Consider the case when $\Omega$ is the unit ball and $\eta(x)=$ $(x, 0)$ for $x \in \partial \Omega$. If $H$ satisfies (18), then a conformal representation of a sphere cap of radius $r=\sqrt{n^{n}} /|H| \geq 1$ with $u \mid \partial \Omega=\eta$ is a solution to (4). If $|H|>\sqrt{n^{n}}$, it can be shown that (4) has no solution.

Proof of Theorem5: Note that the equation in (5) is the Euler-Lagrange equation of the functional $I$, defined by

$$
\begin{equation*}
I(u)=\int_{\Omega}|\nabla u|^{n}+\frac{n H}{n+1} u \cdot u_{1} \wedge \ldots \wedge u_{n} \tag{19}
\end{equation*}
$$

without constraint. Since $I$ is neither bounded from above, nor from below, it has no global maximum nor minimum. We will find a local minimum of $I$ by minimizing $I$ on the subset

$$
M=\left\{u \in W^{1, n}\left(\Omega, R^{n+1}\right): u=\eta \text { on } \partial \Omega,\|u\|_{\infty}|H| \leq \sqrt{n^{n}} \frac{2 n+1}{2 n}\right\}
$$

It is easy to see that $M$ is weakly closed and convex subset of $W^{1, n}\left(\Omega, R^{n+1}\right)$. For any $u \in M$, it follows from (9) that

$$
\begin{gather*}
I(u) \geq \int_{\Omega}|\nabla u|^{n}-\frac{n|H|| | u \mid \|_{\infty}}{(n+1) \sqrt{n^{n}}} \int_{\Omega}|\nabla u|^{n}  \tag{20}\\
\geq \frac{1}{2 n+2} \int_{\Omega}|\nabla u|^{n}
\end{gather*}
$$

So $I$ is coercive. From [23] or [8], $I$ is quasiconvex. By the Theorem II. 4 in [1], $I$ is weakly lower semicontinuous. It follows from the direct method that $I$ has a minimum $u$ in $M$.

We now show that $\|u\|_{\infty} \leq\|\eta\|_{\infty}$. Suppose $k$ is any number satisfying

$$
\begin{equation*}
\|\eta\|_{\infty}|H|<k|H| \leq \sqrt{n^{n}} \frac{2 n+1}{2 n} \tag{21}
\end{equation*}
$$

Let $\phi=\max \{|u|-k, 0\}$. Then $\phi \in W_{0}^{1, n}\left(\Omega, R^{+}\right) \cap L^{\infty}$, and $u-t \phi u \in M$ for sufficiently small $t \geq 0$. It follows from the minimality of $u$,

$$
\begin{aligned}
0 & \geq-\left.\frac{d}{d t}\right|_{t=0} I(u-t \phi u)=\int_{\Omega}<\phi u, D I(u)> \\
& =n \int_{\Omega}|\nabla u|^{n-2} \nabla u \nabla(\phi u)+\frac{|H|}{n+1}(\phi u) \cdot u_{1} \wedge \ldots \wedge u_{n} \\
& \geq n \int_{\Omega}\left(|\nabla u|^{n}-\frac{\left.|H|| | u\right|_{\infty}}{n+1}\left|u_{1} \wedge \ldots \wedge u_{n}\right|\right) \phi+n \int_{\Omega}|\nabla u|^{n-2} \nabla u \cdot u \nabla \phi \\
& \geq \frac{n}{2 n+2} \int_{\{|u|>k\}}|\nabla u|^{n} \phi+n \int_{\{|u|>k\}}|\nabla u|^{n-2}(\nabla u \cdot u)^{2}|u|^{-1} .
\end{aligned}
$$

It follows that $\nabla u=0$ a.e. on $\{|u|>k\}$, which implies that $\nabla \phi=0$ a.e. $\Omega$. So $\phi \equiv 0$, or $|u| \leq k$. As $k$ in (21) is arbitrary, $\|u\|_{\infty} \leq\|\eta\|_{\infty}$, which implies that $\|u\|_{\infty}|H| \leq\|u\|_{\infty}|H|<\sqrt{n^{n}} \frac{2 n+1}{2 n}$. So $u$ is an interior minimum point of $M$ in the norm $\|\cdot\|_{\infty}$; it is then has to be a critical point of $I$ and satisfies (5).

## 3 The Existence of Large Solutions

In Section 2, we showed that if $\|\eta\|_{\infty}|H| \leq \sqrt{n^{n}}$, then the Dirichlet problem (22) has a solution. In this section, we will prove that there is at least another big solution if $\eta$ is small enough. When $n=2$, the existence of multiple solutions of (22) was established in [6][29] under the optimal assumption $0 \neq$ $\|\eta\|_{\infty}|H|<2$. The optimal condition for our case is expected to be $0 \neq$ $\|\eta\|_{\infty}|H|<\sqrt{n^{n}}$, though our proof of Theorem (12) does not yield such an estimate.

Denote by $u_{0}$ the solution we found in Theorem 5 of Section 3. We will solve the problem

$$
\begin{equation*}
\operatorname{div}\left(|\nabla u|^{n-2} \nabla u\right)=H u_{1} \wedge \cdots \wedge u_{n}, u=\eta \text { on } \partial \Omega, \tag{22}
\end{equation*}
$$

for $u=u_{0}+v$ with some $v \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right), v \neq 0$. Note that (22) is the Euler-Lagrange equation of the functional

$$
\begin{equation*}
E(u)=\int_{\Omega}|\nabla u|^{n}+\frac{n H}{n+1} Q(u), \tag{23}
\end{equation*}
$$

without constraint, where $Q(u)=\int_{\Omega} u \cdot u_{1} \wedge \cdots \wedge u_{n}=(n+1) V(u)$. The method is to find a critical point of (23) . . We need some preparations.

Proposition 8 For $a, b \in R^{k},(k \geq 2$ an integer $)$, there holds

$$
\begin{equation*}
|a+b|^{n}=|a|^{n}+|b|^{n}+n|a|^{n-2} a \cdot b+M(a, b) \tag{24}
\end{equation*}
$$

where $M(a, b)$ satisfies

$$
\begin{equation*}
|M(a, b)| \leq n(n-2)(|a|+|b|)^{n-3}|a||b|^{2} . \tag{25}
\end{equation*}
$$

Proof: By the fundamental theorem of calculus,

$$
\begin{align*}
M(a, b) & \equiv|a+b|^{n}-\left(|a|^{n}+|b|^{n}+n|a|^{n-2} a \cdot b\right) \\
& =\int_{0}^{1} \frac{d}{d t}|a+t b|^{n} d t-\left(|b|^{n}+n|a|^{n-2} a \cdot b\right) \\
& =n \int_{0}^{1}|a+t b|^{n-2}\left(a \cdot b+t|b|^{2}\right) d t-\left(|b|^{n}+n|a|^{n-2} a \cdot b\right) \\
& =n \int_{0}^{1} \int_{0}^{t} \frac{d}{d s}|a+s b|^{n-2} a \cdot b d s d t+n \int_{0}^{1} \int_{0}^{1} t \frac{d}{d s}|s a+t b|^{n-2}|b|^{2} d s d t . \tag{26}
\end{align*}
$$

(25) follows from the following estimate: For any $p \geq 1$,

$$
\begin{equation*}
\left.\sup _{0 \leq t \leq 1}\left|\frac{d}{d t}\right| a+\left.t b\right|^{p}\left|\leq p(|a|+|b|)^{p-1}\right| b \right\rvert\, . \tag{27}
\end{equation*}
$$

## Proposition 9

$$
\begin{equation*}
Q\left(u_{0}+v\right)=Q\left(u_{0}\right)+Q(v)+\sum_{i=1}^{n-1} Q_{i}(v) \tag{28}
\end{equation*}
$$

where $Q_{i}(v)$ is homogeneous in $v$ of degree $i$ and homogeneous in $u_{0}$ of degree $n+1-i$.

Proof: Let $g(t)=Q\left(u_{0}+t v\right)$. Then (28) is the Taylor expansion of $g$ at $t=1$, where $Q_{i}(v)=g^{(i)}(0) / i$ ! and $Q(v)=Q_{n+1}(v)$.

Proposition 10

$$
\begin{equation*}
n \int_{\Omega}\left|\nabla u_{0}\right|^{n-2} \nabla u_{0} \nabla v+\frac{n H}{n+1} Q_{1}(v)=0 . \tag{29}
\end{equation*}
$$

Proof: The is just the weak form of the equation (22) ; v serves as a test function.

It follows from (23)-(29) that

$$
\begin{align*}
E\left(u_{0}+v\right)= & \int_{\Omega}\left|\nabla u_{0}\right|^{n}+\frac{n H}{n+1} Q\left(u_{0}\right)+\int_{\Omega}|\nabla v|^{n}  \tag{30}\\
& +\frac{n H}{n+1} Q_{n}(v)+E_{2}(v)+\frac{n H}{n+1} Q(v),
\end{align*}
$$

where

$$
\begin{gather*}
Q_{n}(v)=(n+1) \int_{\Omega} u_{0} \cdot v_{1} \wedge \cdots \wedge v_{n}, \text { by }(17), \\
E_{2}(v)=\int_{\Omega} M\left(\nabla u_{0}, \nabla v\right)+\sum_{i=2}^{n-1} Q_{i}(v) . \tag{31}
\end{gather*}
$$

Since the first two terms of (30) are constant, we are led to the functional

$$
\begin{equation*}
\Phi(v) \equiv \int_{\Omega}|\nabla v|^{n}+\frac{n H}{n+1} Q_{n}(v)+E_{2}(v)+\frac{n H}{n+1} Q(v) . \tag{32}
\end{equation*}
$$

We look at each term in (32). Note that by (9),

$$
\begin{equation*}
\left|Q_{n}(v)\right| \leq C \sup \left|u_{0}\right| \int_{\Omega}|\nabla v|^{n}, \text { where } C=\frac{n+1}{\sqrt{n^{n}}} . \tag{33}
\end{equation*}
$$

The isoperimetric inequality for mappings [2][Theorem 12] implies that if $v \in W_{0}^{1 . n}\left(\Omega, R^{n+1}\right)$ then

$$
\begin{equation*}
|V(v)| \leq \frac{1}{C} A(v)^{\frac{n+1}{n}} \tag{34}
\end{equation*}
$$

where $C=(n+1) \omega_{n}^{\frac{1}{n}}$ and $\omega_{n}$ is the area of the unit $n$-sphere $S^{n}$. In terms of $Q(v)=(n+1) V(v)$ and $\int_{\Omega}|\nabla v|^{n}$, it follows from (17) that

$$
\begin{equation*}
|Q(v)|^{\frac{n}{n+1}} \leq \frac{1}{S} \int_{\Omega}|\nabla v|^{n}, \text { where } S=n^{\frac{n}{2}} \omega_{n}^{\frac{1}{n+1}} \tag{35}
\end{equation*}
$$

To estimate $E_{2}(v)$, we first notice that $Q(u)=R(u, u)$ and

$$
\begin{align*}
\sum_{i=2}^{n-1} Q_{i}(v) & =Q\left(u_{0}+v\right)-Q\left(u_{0}\right)-Q_{1}(v)-Q_{n}(v)-Q(v) \\
& =\left[R\left(u_{0}, u_{0}+v\right)-R\left(u_{0}, u_{0}\right)-R\left(u_{0}, v\right)\right]+  \tag{36}\\
& {\left[R\left(v, u_{0}+v\right)-R(v, v)-n R\left(u_{0}, v\right)\right]-(n+1) R\left(v, u_{0}\right) }
\end{align*}
$$

By (14), we have

$$
\begin{align*}
& \left|R\left(v, v+u_{0}\right)-R(v, v)-n R\left(u_{0}, v\right)\right| \\
& =\left|\int_{\Omega} \int_{0}^{1}\left[\int_{0}^{t} \frac{d}{d s} \sum_{i=0}^{n} u_{0} \cdot\left(v+s u_{0}\right)_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge\left(v+s u_{0}\right)_{n} d s\right] d t\right| \\
& \leq C| | u_{0}\left\|_{\infty}\right\| \nabla u_{0}\left\|_{L^{n}}\right\| \nabla v\left\|_{L^{n}}| ||\nabla v|+\mid \nabla u_{0}\right\| \|_{L^{n}}^{n-2} . \tag{37}
\end{align*}
$$

$\left|R\left(u_{0}, u_{0}+v\right)-R\left(u_{0}, u_{0}\right)-R\left(u_{0}, v\right)\right|$
$=\left|\int_{\Omega} \int_{0}^{1} \sum_{i=0}^{n} u_{0} \cdot\left(s u_{0}+t v\right)_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge\left(s u_{0}+t v\right)_{n}\right|_{s=0}^{s=1} d t \mid$
$=\left|\int_{\Omega} \int_{0}^{1} \int_{0}^{1} \frac{d}{d s} \sum_{i=0}^{n} u_{0} \cdot\left(s u_{0}+t v\right)_{1} \wedge \ldots \wedge v_{i} \wedge \ldots \wedge\left(s u_{0}+t v\right)_{n} d s d t\right|$
$\leq C\left\|u_{0}\right\|_{\infty}\left\|\nabla u_{0}\right\|_{L^{n}}\|\nabla v\|_{L^{n}}\left\||\nabla v|+\left|\nabla u_{0}\right|\right\|_{L^{n}}^{n-2}$.

$$
\begin{equation*}
\left|R\left(v, u_{0}\right)\right| \leq C\|\nabla v\|_{L^{n}}\left\|\nabla u_{0}\right\|_{L^{n}}^{n} . \tag{39}
\end{equation*}
$$

By (31), (25) and (36)-(39), we get

$$
\begin{align*}
\left|E_{2}(v)\right| \leq & \int_{\Omega} n(n-2)\left(\left|\nabla u_{0}\right|+|\nabla v|\right)^{n-3}\left|\nabla u_{0}\right||\nabla v|^{2}+\left|\sum_{i=2}^{n-1} Q_{i}(v)\right| \\
\leq & C \int_{\Omega}\left(\left|\nabla u_{0}\right|^{n-2}|\nabla v|^{2}+\left|\nabla u_{0}\right|^{2}|\nabla v|^{n-2}\right)+ \\
& C \int_{\Omega} \sum_{i=2}^{n-1}\left|u_{0}\right|_{\infty}\left|\nabla u_{0}\right||\nabla v|\left(|\nabla v|^{n-2}+\left|\nabla u_{0}\right|^{n-2}\right)+C|\nabla v|\left|\nabla u_{0}\right|^{n} \\
\leq & C_{0} \int_{\Omega}^{n-1}\left|\nabla u_{i=1}^{n-1}\right|^{n-i}|\nabla v|^{i}+C|\nabla v|\left|\nabla u_{0}\right|^{n} . \tag{40}
\end{align*}
$$

Note that $\Phi$ is unbounded from above and below, and it is a typical case not satisfying the Palais-Smale conditions. The standard variational method
fails to give the existence of a critical point. In the case $n=2$, where $E_{2}$ does not appear in $\Phi$, Brezis and Coron [6] was able to find a nontrivial critical point of $\Phi$ as a proper dilation of a minimum of $\int_{\Omega}|\nabla v|^{2}+\frac{2 H}{3} Q_{2}(v)$ subject to $Q(v)=$ constant. For $n \geq 3$, the terms of $\Phi$ have at least three different homogeneities, therefore, the method in [6] is unlikely to work. Our method is to apply a mountain pass theorem of Ambrosetti-Rabinowitz [3] in a min-max scheme. We will use the following form of the theorem in [3], as used by Brezis and Nirenberg[7] in solving elliptic equations with critical exponents.

Theorem 11 [3][7]Assumption: Let $\Phi$ be a $C^{1}$ function on a Banach space $E$. Suppose there exists a neighborhood $U$ of 0 in $E$ and a constant $\rho$ such that $\Phi(u) \geq \rho$ for every $u \in \partial U$, and

$$
\Phi(0)<\rho \text { and } \Phi(v)<\rho \text { for some } v \notin U \text {. }
$$

Set $c=\inf _{p \in P} \max _{w \in p} \Phi(w) \geq \rho$, where $P$ denotes the class of paths joining 0 to $v$.

Conclusion: There is a sequence $\left\{u_{i}\right\}$ in $E$ such that $\Phi\left(u_{i}\right) \rightarrow c$ and

$$
\Phi^{\prime}\left(u_{i}\right) \rightarrow 0 \text { in } E^{*} .
$$

The advantage of this theorem is that it does not require (PS)-condition. We will show that a subsequence of $\left\{u_{i}\right\}$ converges to a nontrivial critical point of $\Phi$. Our result is stated as follows.

Theorem $12 \eta \in W^{1, n}\left(\Omega, R^{n+1}\right)$ and $\|\eta\|_{\infty}+\|\nabla \eta\|_{L^{n}(\partial \Omega)}$ is small enough, then the problem (22) has at least two solutions.

Remark 13 One solution is the small solution $u_{0}$ found in Section 2; it satisfies $\left\|u_{0}\right\|_{\infty} \leq\|\eta\|_{\infty}$ and is a minimum of $E$ in $M$. Thus

$$
\frac{1}{2 n+2} \int_{\Omega}\left|\nabla u_{0}\right|^{n} \leq E\left(u_{0}\right) \leq E(\bar{\eta}),
$$

where $\bar{\eta}(x)=|x| \eta\left(\frac{x}{|x|}\right)$ is a special extension of $\eta$. Thus

$$
\int_{\Omega}\left|\nabla u_{0}\right|^{n} \leq \int_{\Omega}|\nabla \bar{\eta}|^{n}+2 n H Q(\bar{\eta}) \leq C_{0} \int_{\Omega}|\nabla \bar{\eta}|^{n} \leq C_{1}\left(\|\eta\|_{\infty}+\|\nabla \eta\|_{L^{n}(\partial \Omega)}\right) .
$$

It follows that $\|\eta\|_{\infty}+\|\nabla \eta\|_{L^{n}(\partial \Omega)}$ is small implies that $\left\|u_{0}\right\|_{\infty}+\left\|\nabla u_{0}\right\|_{L^{n}(\Omega)}$ is also small. The smallness condition used in the proof actually is referred to $u_{0}$.

We now start the proof of Theorem 12 with verifying the conditions in Theorem 11.

Proposition 14 There are numbers $\delta, \rho>0$ such that

$$
\Phi(v) \geq \rho \text { for } v \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right) \text { with }\|\nabla v\|_{L^{n}(\Omega)}=\delta \text {. }
$$

Proof: $\mathrm{By}(32),(33),(35),(40)$ and the Hölder inequality, for any $\epsilon>0$, there are numbers $C_{0}, C(\epsilon)$, such that

$$
\begin{array}{r}
\Phi(v) \geq \int_{\Omega}|\nabla v|^{n}-\left.C_{0}\left|\left\|u_{0}\right\|_{\infty} \int_{\Omega}\right| \nabla v\right|^{n}-\epsilon \int_{\Omega}|\nabla v|^{n}- \\
C(\epsilon)\left(\left|u_{0}\right|_{\infty}+\left\|\nabla u_{0}\right\|_{L^{n}(\Omega)}\right)-C_{0}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{n+1}{n}} .
\end{array}
$$

Fix $\epsilon=\frac{1}{4}$ and a number $\delta>0$ such that $C_{0} \delta \leq \frac{1}{8}$. Suppose $u_{0}$ satisfies $C_{0}\left\|u_{0}\right\|_{\infty} \leq \frac{1}{4}$ and $C\left(\frac{1}{4}\right)\left(\left|u_{0}\right|_{\infty}+\left\|\nabla u_{0}\right\|_{L^{n}(\Omega)}\right) \leq \frac{1}{16} \delta^{n}$. Then for any $v \in$ $W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$ with $\|\nabla v\|_{L^{n}(\Omega)}=\delta$, we have

$$
\begin{gathered}
\Phi(v) \geq \frac{1}{4} \int_{\Omega}|\nabla v|^{n}-C_{0}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{n+1}{n}}-C(\epsilon)\left(\left|u_{0}\right|_{\infty}+\left\|\nabla u_{0}\right\|_{L^{n}(\Omega)}\right) \\
\geq \frac{1}{8} \delta^{n}-C(\epsilon)\left(\left|u_{0}\right|_{\infty}+\left\|\nabla u_{0}\right\|_{L^{n}(\Omega)}\right) \geq \frac{1}{16} \delta^{n} .
\end{gathered}
$$

The proposition holds with $\rho=\frac{1}{16} \delta^{n}$.
Proposition 15 There is a $v \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$ such that

$$
\begin{gather*}
\Phi(v) \leq 0 \\
\sup _{0 \leq t} \Phi(t v)<\frac{S^{n+1}}{|H|^{n}(n+1)} . \tag{41}
\end{gather*}
$$

The proof of Proposition 15 will be given later. Now we prove Theorem 12.

Proof of Theorem 12: By the theorem of Ambrosetti-Rabinowitz above and Propositions 14 and 15 , there exists $\left\{v^{i}\right\} \subset W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$ such that as $i \rightarrow \infty$,

$$
\begin{gather*}
\Phi\left(v^{i}\right)=\int_{\Omega}\left|\nabla v^{i}\right|^{n}+\frac{n H}{n+1} Q_{n}\left(v^{i}\right)+  \tag{42}\\
E_{2}\left(v^{i}\right)+\frac{n H}{n+1} Q\left(v^{i}\right) \xrightarrow{\rightarrow},
\end{gather*}
$$

where

$$
\begin{equation*}
c=\inf _{P} \max _{v \in P}\{\Phi(v)\}, \tag{43}
\end{equation*}
$$

and

$$
\begin{align*}
\frac{1}{n} \Phi^{\prime}\left(v^{i}\right) & =-\operatorname{div}\left(\left|\nabla v^{i}\right|^{n-2} \nabla v^{i}\right)+\frac{H}{n+1} Q_{n}^{\prime}\left(v^{i}\right)+  \tag{44}\\
& \frac{1}{n} E_{2}^{\prime}\left(v^{i}\right)+H v_{1}^{i} \wedge \cdots \wedge v_{n}^{i} \rightarrow 0 \text { in } W^{-1, n^{\prime}}
\end{align*}
$$

where $n^{\prime}=\frac{n}{n-1}$. Multiply (44) by $v^{i}$ and integrate. We get

$$
\begin{equation*}
\int_{\Omega}\left|\nabla v^{i}\right|^{n}+\frac{H}{n+1}<Q_{n}^{\prime}\left(v^{i}\right), v^{i}>+\frac{1}{n}<E_{2}^{\prime}\left(v^{i}\right), v^{i}>+H Q\left(v^{i}\right) \rightarrow 0 . \tag{45}
\end{equation*}
$$

We claim that

$$
\begin{equation*}
\int_{\Omega}\left|\nabla v^{i}\right|^{n} \leq C \tag{46}
\end{equation*}
$$

for some constant $C$. To prove (46), we first note that since $Q_{n}\left(v^{i}\right)$ is homogeneous in $v^{i}$ of degree $n$,

$$
\begin{equation*}
\left|<Q_{n}^{\prime}\left(v^{i}\right), v^{i}>\left|=\left|n Q_{n}\left(v^{i}\right)\right| \leq \frac{n(n+1)}{\sqrt{n^{n}}}\left\|u_{0}\right\|_{\infty} \int_{\Omega}\right| \nabla v^{i}\right|^{n} ; \tag{47}
\end{equation*}
$$

and by (40) and the Hölder inequality, for $\epsilon>0$, there is a constant $C(\epsilon)$, such that

$$
\begin{gather*}
\left|E_{2}\left(v^{i}\right)\right| \leq C(\epsilon) \int_{\Omega}\left|\nabla u_{0}\right|^{n}+\epsilon \int_{\Omega}\left|\nabla v^{i}\right|^{n}  \tag{48}\\
\left|<E_{2}^{\prime}\left(v^{i}\right), v^{i}>\left|\leq C(\epsilon) \int_{\Omega}\right| \nabla u_{0}\right|^{n}+\epsilon \int_{\Omega}\left|\nabla v^{i}\right|^{n} . \tag{49}
\end{gather*}
$$

Now look at the difference of (42) and (45), we then get

$$
\frac{H}{n+1} Q_{n}\left(v^{i}\right)+\frac{H}{n+1} Q\left(v^{i}\right)+\frac{1}{n}<E_{2}^{\prime}\left(v^{i}\right), v^{i}>-E_{2}\left(v^{i}\right) \rightarrow-c .
$$

It follows for some constant $C$, depending on $\epsilon$,

$$
\begin{equation*}
\left|Q\left(v^{i}\right)\right| \leq \epsilon \int_{\Omega}\left|\nabla v^{i}\right|^{n}+C(\epsilon) . \tag{50}
\end{equation*}
$$

Combining (50) with (42), we get (46) . As in [18], we may assume, by passing to a subsequence, that $v^{i}$ weakly converges to a $v$ in $W^{1, n}\left(\Omega, R^{n+1}\right)$, and strongly converges to $v$ in $W^{1, p}\left(\Omega, R^{n+1}\right)$ for any $p \in[1, n)$.

We claim that $v$ is nontrivial. For otherwise, $v \equiv 0$ implies that

$$
\begin{align*}
<Q_{n}^{\prime}\left(v^{i}\right), v^{i}> & =n Q_{n}\left(v^{i}\right)=(n+1) \int_{\Omega} u_{0} \cdot v_{1}^{i} \wedge \cdots \wedge v_{n}^{i}, \rightarrow 0  \tag{51}\\
& E_{2}\left(v^{i}\right) \rightarrow 0,<E_{2}^{\prime}\left(v^{i}\right), v^{i}>\rightarrow 0 .
\end{align*}
$$

By passing to a subsequence if necessary, we may assume further that $\int_{\Omega}\left|\nabla v^{i}\right|^{n} \rightarrow$ $l$. It follows that $Q\left(v^{i}\right) \rightarrow-\frac{l}{H}$ by (45). By (42), we have

$$
\begin{equation*}
l+\frac{n H}{n+1}\left(-\frac{1}{H}\right) l \rightarrow c . \tag{52}
\end{equation*}
$$

It follows that $c=\frac{l}{n+1}$. On the other hand, by isoperimetric inequality,

$$
l \geq S\left|\frac{l}{H}\right|^{\frac{n}{n+1}}
$$

which implies $l \geq \frac{S^{n+1}}{|H|^{n}}$. Therefore,

$$
c \geq \frac{S^{n+1}}{|H|^{n}(n+1)}
$$

This is a contradiction, because Proposition 15 implies that $c<\frac{S^{n+1}}{H^{n}(n+1)}$. So $v$ is nontrivial. Taking the limit in (44), we have that $v$ satisfies $\Phi^{\prime}(v)=0$, or equivalently, $u=u_{0}+v$ is a solution.

The rest of this section is devoted to the proof of Proposition 15. The case $n=2$ has been shown in [6]. We generalize the argument in [6] to higher dimensions.

For $v \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right)$, denote

$$
\begin{gather*}
E_{3}(v)=\int_{\Omega}|\nabla v|^{n}+\frac{n H}{n+1} Q_{n}(v)  \tag{53}\\
R(v)=\frac{E_{3}(v)}{|Q(v)|^{\frac{n}{n+1}}} ;  \tag{54}\\
S=\inf \left\{\frac{\int_{\Omega}|\nabla v|^{n}}{|Q(v)|^{\frac{n}{n+1}}}, Q(v) \neq 0, v \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right)\right\} . \tag{55}
\end{gather*}
$$

Define

$$
\begin{equation*}
J=\inf \left\{T(v): Q(v) \neq 0, v \in W_{0}^{1, n}\left(\Omega, R^{n+1}\right)\right\} \tag{56}
\end{equation*}
$$

We first prove

## Proposition $16 \quad J<S$.

Proof: Suppose $0 \in \Omega$ and $\nabla u(0) \neq 0$. Choose a coordinate basis $e_{1, \ldots}, e_{n+1}$ for $R^{n+1}$ that has the same orientation as the canonical basis of $R^{n+1}$ such that

$$
\begin{equation*}
\gamma \equiv \frac{\partial u}{\partial x_{1}}(0) \cdot e_{1}+\ldots+\frac{\partial u}{\partial x_{n}}(0) \cdot e_{n}<0 . \tag{57}
\end{equation*}
$$

Let $v: R^{n} \rightarrow S^{n}$ be the stereographic projection:

$$
\begin{equation*}
v(x)=\frac{(2 x,-2)}{1+|x|^{2}}, x \in R^{n} \tag{58}
\end{equation*}
$$

( $v$ is written in the coordinate $e_{1, \cdots}, e_{n+1}$ ). For $\epsilon>0$, consider the map

$$
v^{\epsilon}(x)=\frac{(2 x,-2 \epsilon)}{\epsilon^{2}+|x|^{2}}
$$

Let $R>0$ be a number such that $B_{4 R} \equiv B_{4 R}(0) \subseteq \Omega$. Let $\xi \in C_{0}^{1}\left(B_{2 R},[0,1]\right)$ be a cut-off function such that $\xi=1$ on $B_{R}$. Note that $\xi v^{\epsilon} \in C_{0}^{1}\left(\Omega, R^{n+1}\right)$ and the following properties of $v^{\epsilon}$ can be easily verified:

$$
v^{\epsilon}(x)=\frac{1}{\epsilon} v\left(\frac{x}{\epsilon}\right),
$$

$$
\begin{align*}
& \left|v^{\epsilon}(x)\right|=\frac{2}{\sqrt{\epsilon^{2}+|x|^{2}}}  \tag{59}\\
& \left|\nabla v^{\epsilon}(x)\right| \leq \frac{C}{\epsilon^{2}+|x|^{2}}
\end{align*}
$$

for a constant $C$ independent of $\epsilon$ and $x$.
We shall establish

$$
\begin{equation*}
T\left(\xi v^{\epsilon}\right)=S+c_{0} \epsilon+O\left(\epsilon^{1+\alpha}\right) \text { as } \epsilon \rightarrow 0 \tag{60}
\end{equation*}
$$

where $c_{0}<0$ and $\alpha \in(0,1)$ are constants. Here, as a notation, $O(f)$ denotes a quantity satisfying $|O(f)| \leq C|f|$ for some constant $C$. The inequality of the Proposition 16 follows by taking $\epsilon$ small enough.

We now proceed to show (60). By the mean value theorem,

$$
\begin{equation*}
|f(a+b)-f(a)|=O\left(\sup _{0 \leq t \leq 1}\left|f^{\prime}(a+t b)\right|\right)|b| \tag{61}
\end{equation*}
$$

Applying this to $f(a)=|a|^{n}$ with $a=\xi \nabla v^{\epsilon}, b=\nabla \xi v^{\epsilon}$, we have

$$
\begin{align*}
& \int_{\Omega}\left|\nabla\left(\xi v^{\epsilon}\right)\right|=\int_{R^{n}}\left|\xi \nabla v^{\epsilon}+\nabla \xi v^{\epsilon}\right|^{n}  \tag{62}\\
& =\int_{R^{n}}\left|\xi \nabla v^{\epsilon}\right|^{n}+O\left(\int_{R^{n}}\left(\left|\xi \nabla v^{\epsilon}\right|+\left|\nabla \xi v^{\epsilon}\right|\right)^{n-1}\left|\nabla \xi v^{\epsilon}\right|\right) .
\end{align*}
$$

Since $v^{\epsilon}$ is conformal and $v^{\epsilon}\left(R^{n}\right)$ is a sphere of radius $\frac{1}{\epsilon}$, we have

$$
\begin{equation*}
\int_{R^{n}}\left|\nabla v^{\epsilon}\right|^{n}=\sqrt{n^{n}} \cdot \operatorname{area}\left(v^{\epsilon}\left(R^{n}\right)\right)=\frac{\sqrt{n^{n}} \omega_{n}}{\epsilon^{n}} \tag{63}
\end{equation*}
$$

On the other hand, by (59),

$$
\begin{equation*}
\int_{R^{n}}\left(\xi^{n}-1\right)\left|\nabla v^{\epsilon}\right|^{n}=O\left(\int_{|x| \geq R}\left|\nabla v^{\epsilon}\right|^{n}\right)=O\left(\int_{R}^{\infty} \frac{r^{n-1}}{r^{2 n}} d r\right)=O(1) . \tag{64}
\end{equation*}
$$

Similarly,

$$
\begin{gather*}
O\left(\int_{R^{n}}\left(\left|\xi \nabla v^{\epsilon}\right|\right)^{n-1}\left|\nabla \xi v^{\epsilon}\right|\right)=O(1)  \tag{65}\\
O\left(\int_{R^{n}}\left|\nabla \xi v^{\epsilon}\right|^{n}\right)=O(1)
\end{gather*}
$$

It follows from (62)-(65)

$$
\begin{equation*}
\int_{\Omega}\left|\nabla\left(\xi v^{\epsilon}\right)\right|=\frac{\sqrt{n^{n}} \omega_{n}}{\epsilon^{n}}+O(1) \tag{66}
\end{equation*}
$$

We now estimate $Q\left(\xi v^{\epsilon}\right)$. Applying (61) to $f(a)=v_{1} \wedge \cdots \wedge v_{n}$ (where $\left.a=\left(v_{j}^{i}\right)\right)$ with $a=\xi \nabla v^{\epsilon}, b=\nabla \xi v^{\epsilon}$, we have

$$
\begin{align*}
Q\left(\xi v^{\epsilon}\right) & =\int_{\Omega} \xi v^{\epsilon} \cdot\left(\xi v^{\epsilon}\right)_{1} \wedge \cdots \wedge\left(\xi v^{\epsilon}\right)_{n} \\
& =\int_{\Omega} \xi^{n+1} v^{\epsilon} \cdot v_{1}^{\epsilon} \wedge \cdots \wedge v_{n}^{\epsilon}+O\left(\int_{R^{n}}\left|\xi v^{\epsilon}\right|\left(\left|\xi \nabla v^{\epsilon}\right|+\left|\nabla \xi v^{\epsilon}\right|\right)^{n-1}\left|\nabla \xi v^{\epsilon}\right|\right) \tag{67}
\end{align*}
$$

Recall that $Q\left(v^{\epsilon}\right) /(n+1)$ is the oriented volume of $v^{\epsilon}\left(R^{n}\right)$. So we have

$$
\begin{equation*}
Q\left(v^{\epsilon}\right)= \pm(n+1) \operatorname{vol}\left(v^{\epsilon}\left(R^{n}\right)\right)= \pm \frac{\omega_{n}}{\epsilon^{n+1}} \tag{68}
\end{equation*}
$$

Similarly to (64) and (65), we have

$$
\begin{gather*}
\int_{\Omega}\left(\xi^{n+1}-1\right) v^{\epsilon} \cdot v_{1}^{\epsilon} \wedge \cdots \wedge v_{n}^{\epsilon}=O(1)  \tag{69}\\
O\left(\int_{R^{n}}\left|\xi v^{\epsilon}\right|\left(\left|\xi \nabla v^{\epsilon}\right|+\left|\nabla \xi v^{\epsilon}\right|\right)^{n-1}\left|\nabla \xi v^{\epsilon}\right|\right)=O(1) \tag{70}
\end{gather*}
$$

So (67) - (70) imply that

$$
\begin{equation*}
\left|Q\left(v^{\epsilon}\right)\right|=\frac{\omega_{n}}{\epsilon^{n+1}}+O(1) \tag{71}
\end{equation*}
$$

Similar argument applies to $Q\left(\xi v^{\epsilon}\right)$, and we have

$$
\begin{align*}
& \frac{1}{n+1} Q_{n}\left(\xi v^{\epsilon}\right)=\int_{\Omega} u \cdot\left(\xi v^{\epsilon}\right)_{1} \wedge \cdots \wedge\left(\xi v^{\epsilon}\right)_{n} \\
& =\int_{\Omega} \xi^{n} u \cdot v_{1}^{\epsilon} \wedge \cdots \wedge v_{n}^{\epsilon}+O\left(\int_{R^{n}}|u|\left(\left|\xi \nabla v^{\epsilon}\right|+\left|\nabla \xi v^{\epsilon}\right|\right)^{n-1}\left|\nabla \xi v^{\epsilon}\right|\right)  \tag{72}\\
& =\int_{\Omega} \xi^{n} u \cdot v_{1}^{\epsilon} \wedge \cdots \wedge v_{n}^{\epsilon}+O(1)
\end{align*}
$$

Denote $\tilde{u}=\xi^{n} u$. Since $\tilde{u}$ is in $C_{0}^{1, \alpha}\left(B_{2 R}\right)$ by the regularity theorem in [24], we have that for all $x \in R^{n}$,

$$
\begin{align*}
& \tilde{u}(x)=\tilde{u}(0)+\nabla \tilde{u}(0) x+O\left(\|\nabla \tilde{u}\|_{C^{\alpha}}|x|^{1+\alpha}\right) \\
& =u(0)+\nabla u(0) x+O\left(\|u\|_{C^{1, \alpha}\left(B_{2 R}\right)}|x|^{1+\alpha}\right) . \tag{73}
\end{align*}
$$

Therefore, by conformal invariance of $Q_{n}$ and (73),

$$
\begin{align*}
\int_{\Omega} \tilde{u} \cdot v_{1}^{\epsilon} \wedge \cdots \wedge v_{n}^{\epsilon} & =\frac{1}{\epsilon^{n}} \int_{R^{n}} \tilde{u}(\epsilon x) \cdot v_{1} \wedge \cdots \wedge v_{n} d x \\
& =\frac{1}{\epsilon^{n}}\left(\int_{R^{n}}(u(0)+\epsilon \nabla u(0) x) \cdot v_{1} \wedge \cdots \wedge v_{n} d x\right)+ \\
& O\left(\epsilon^{1-n+\alpha}| | \nabla u \|_{C^{\alpha}\left(B_{2 R}\right)} \int_{|x| \leq \frac{2 R}{\epsilon}}|x|^{1+\alpha}|\nabla v|^{n}\right) . \tag{74}
\end{align*}
$$

We have

$$
\begin{equation*}
\int_{R^{n}} u(0) \cdot v_{1} \wedge \cdots \wedge v_{n} d x=0 \tag{75}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{|x| \leq \frac{2 R}{\epsilon}}|x|^{1+\alpha}|\nabla v|^{n}=O(1)+\int_{1 \leq|r| \leq \frac{2 R}{\epsilon}} r^{\alpha-n}=O(1)+O\left(\epsilon^{n-\alpha+1}\right) . \tag{76}
\end{equation*}
$$

Next we show that

$$
\begin{equation*}
\int_{R^{n}} \nabla u(0) x \cdot v_{1} \wedge \cdots \wedge v_{n} d x=c \Sigma_{i=1}^{n} \frac{\partial u}{\partial x_{i}} \cdot e_{i}=c \gamma \tag{77}
\end{equation*}
$$

Proof of (77) : By (17),

$$
\begin{aligned}
& \int_{R^{n}} \nabla u(0) x \cdot v_{1} \wedge \cdots \wedge v_{n} d x \\
& =-\frac{1}{n} \int_{R^{n}} \sum_{i=1}^{n}(\nabla u(0) x)_{i} \cdot v_{1} \wedge \cdots \wedge v_{i}^{v} \wedge \cdots \wedge v_{n} d x \\
& =-\frac{1}{n} \int_{R^{n}} \sum_{i=1}^{n} \frac{\partial u}{\partial x_{i}} \cdot\left[e_{1} \wedge \cdots \wedge(x,-1) \wedge \cdots \wedge e_{n}\right] \frac{1}{\left(1+|x|^{2}\right)^{n}} d x \\
& =-\frac{1}{n} \int_{R^{n}} \sum_{i=1}^{n}\left(\frac{\partial u}{\partial x_{i}} \cdot e_{i}\right) e_{i} \cdot\left[e_{1} \wedge \cdots \wedge\left(-e_{n+1}\right) \wedge \cdots \wedge e_{n}\right] \frac{1}{\left(1+|x|^{2}\right)^{n}} d x \\
& =c^{\prime} \sum_{i=1}^{n} \frac{\partial u}{\partial x_{i}} \cdot e_{i}=c^{\prime} \gamma,
\end{aligned}
$$

where $c^{\prime}=\int_{R^{n}} \frac{d x}{\left(1+|x|^{2}\right)^{n}}$. So (72)-(77) together imply that

$$
\begin{equation*}
\int_{\Omega} u \cdot\left(\xi v^{\epsilon}\right)_{1} \wedge \cdots \wedge\left(\xi v^{\epsilon}\right)_{n}=c^{\prime} \gamma \epsilon^{1-n}+O\left(\epsilon^{1-n+\epsilon}\right) . \tag{78}
\end{equation*}
$$

It follows from (66) (71) (440) that

$$
\begin{aligned}
& T\left(\xi v^{\epsilon}\right)=\left[\epsilon^{-n} \omega_{n}^{\frac{n}{n+1}}+O(1)\right]^{-1}\left(\epsilon^{-n} n^{n / 2} \omega_{n}+O(1)+\frac{n H c^{\prime}}{n+1} \gamma+O\left(\epsilon^{-n+1+\alpha}\right)\right) \\
& =S+c_{0} \epsilon+O\left(\epsilon^{1+\alpha}\right), \text { where } c_{0}=\frac{n H c}{n+1} \gamma<0 .
\end{aligned}
$$

This finishes the proof of (60).
Proof of Proposition 15: Let $\xi v^{\epsilon}$ be as in the proof Proposition 16 such that $T\left(\xi v^{\epsilon}\right)<S$. It is easy to check the $\pm \operatorname{sign}$ in $(68)$ is $(-1)^{n}$. Take $v=\xi v^{\epsilon}$ for $n$ odd. Take $v=-\xi v^{\epsilon}$ for $n$ even; so $T(v)=T\left(\xi v^{\epsilon}\right)$ and $Q(v)=-Q\left(\xi v^{\epsilon}\right)$. Thus for any $n, T(v)<S$ and $Q(v)<0$. We may also assume that $\Phi(v) \leq 0$, by replacing $v$ by $\lambda v$ for large $\lambda>0$. Consider

$$
\begin{equation*}
\Phi^{*}(t v) \equiv E_{3}(t v)+\frac{n H}{n+1} Q(t v)=t^{n} E_{3}(v)+\frac{n H}{n+1} t^{n+1} Q(v) \tag{79}
\end{equation*}
$$

It is easy to check that $\Phi^{*}$ has a maximum at $t=-\frac{E_{3}(v)}{Q(v) H}$, with maximum value

$$
\begin{equation*}
\Phi^{*}(t v)=\left[\frac{E_{3}(v)}{|Q(v)|^{\frac{n}{n+1}}}\right]^{n+1} \frac{1}{|H|^{n}(n+1)}<\frac{S^{n+1}}{|H|^{n}(n+1)} \tag{80}
\end{equation*}
$$

To show (41), we need to assume that $u_{0}$ is small, say, $\int_{\Omega}\left|\nabla u_{0}\right|^{n} \leq 1$, then by (40)

$$
\begin{align*}
\left|E_{2}(v)\right| & \leq C \int_{\Omega} \sum_{i=1}^{n-1}\left|\nabla u_{0}\right|^{n-i}|\nabla v|^{i} \\
& \leq C \sum_{i=1}^{n-1}\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{n-i}{n}}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{i}{n}}  \tag{81}\\
& \leq C\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \sum_{i=1}^{n-1}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{i}{n}}
\end{align*}
$$

It follows

$$
\begin{align*}
\Phi(t v) & \leq t^{n} E_{3}(v)+\frac{n H}{n+1} t^{n+1} Q(v)+C\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \sum_{i=1}^{n-1} t^{i}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{i}{n}} \\
& \equiv \Phi^{* *}(t, v) . \tag{82}
\end{align*}
$$

By the construction of $v$, we have that

$$
\begin{align*}
& \epsilon^{n} E_{3}(v)=\sqrt{n^{n}} \omega_{n}+O(\epsilon) ; \epsilon^{n} \int_{\Omega}|\nabla v|^{n}=\sqrt{n^{n}} \omega_{n}+O(\epsilon) ;  \tag{83}\\
& \epsilon^{n+1} Q(v)=-\omega_{n}+O\left(\epsilon^{n+1}\right) .
\end{align*}
$$

Therefore, there are positive numbers $C_{1}, C_{2}, C_{3}$, such that for any number $\beta>0$,

$$
\begin{align*}
\Phi^{* *}(\beta \epsilon, v) & =\beta^{n} \epsilon^{n} E_{3}(v)+\frac{n H}{n+1} \beta^{n+1} \epsilon^{n+1} Q(v)+ \\
& C\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \sum_{i=1}^{n-1} \beta^{i} \epsilon^{i}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{i}{n}}  \tag{84}\\
& \leq C_{1} \beta^{n}-C_{2} \beta^{n+1}+C_{3}\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \sum_{i=1}^{n-1} \beta^{i} .
\end{align*}
$$

It follows that there is a $\beta^{*}$ such that $\Phi^{* *}(\beta \epsilon, v) \leq 0$ for all $0<\epsilon \ll 1$ and $\beta \geq \beta^{*}$. By (82), we have,

$$
\begin{align*}
\sup _{0 \leq t} \Phi(t v) & \leq \sup _{0 \leq t \leq \beta^{*} \epsilon} \Phi^{* *}(t, v) \\
& \leq \sup _{0 \leq t} \Phi^{*}(t, v)+\sup _{0 \leq t \leq \beta^{*} \epsilon} C\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \sum_{i=1}^{n-1} t^{i}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{i}{n}} \\
& \leq \sup _{0 \leq t} \Phi^{*}(t, v)+C\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \sum_{i=1}^{n-1} \beta^{* i} \epsilon^{i}\left(\int_{\Omega}|\nabla v|^{n}\right)^{\frac{i}{n}} \\
& \leq \sup _{0 \leq t} \Phi^{*}(t, v)+C_{4}\left(\int_{\Omega}\left|\nabla u_{0}\right|^{n}\right)^{\frac{1}{n}} \tag{85}
\end{align*}
$$

where $C_{4}$ depends on $\beta^{*}$. Since $\sup _{0 \leq t} \Phi^{*}(t, v)<\frac{S^{n+1}}{H^{n}(n+1)}$, $\sup _{0 \leq t} \Phi(t v)<$ $\frac{S^{n+1}}{H^{n}(n+1)}$, if $\int_{\Omega}\left|\nabla u_{0}\right|^{n}$ is small enough.

## 4 Regularity of Conformal Solutions

Our result is
Theorem 17 If $u$ is a conformal solution of (1) and $f$ satisfies (2), then $u \in C^{1, \alpha}(\Omega)$ for some $\alpha \in(0,1)$. If $u=\eta$ on $\partial \Omega$ and $\eta \in C^{0}(\partial \Omega)$, then $u \in C^{0}(\bar{\Omega})$.

When $n=2$, this theorem was proved by Grüter in 1980 [16]. We will use the main idea of the proof in [16].

Consider the set $G$ of good points of $u$ defined by

$$
\begin{aligned}
G=\{x \in \Omega: & u \text { is approximately differentiable at } x, \text { and } \\
& \left.x \text { is a Lebesgue point of }|\nabla u|^{n}, \text { and }|\nabla u|(x) \neq 0\right\} .
\end{aligned}
$$

Here $u$ is approximately differentiable at a point $x_{0}$ with approximate differential $\nabla u\left(x_{0}\right)$, by definition, if there is a $u_{0} \in R^{k}$ such that for every $\epsilon>0$,

$$
\Phi^{n}\left[L^{n}\left\lfloor\Omega \backslash\left\{x:\left|u(x)-u_{0}-\nabla u\left(x_{0}\right)\left(x-x_{0}\right)\right| \leq \epsilon\left|x-x_{0}\right|\right\}, x_{0}\right]=0,\right.
$$

where $\Phi^{n}$ denotes the n-dimensional density and $L^{n}\lfloor\Omega$ is the Lebesgue measure, restricted to $\Omega$.

We will need the following property for functions in $W^{1, n}\left(\Omega, R^{k}\right)$.
Proposition 18 ([12] [Theorem 4.5.9]) If $u \in W^{1, n}\left(\Omega, R^{k}\right)$, then $u$ has weak derivative and approximate differential almost everywhere, and when both exist, they coincide.

For a proof, see [12]. Next, we have
Lemma 19 Suppose $u \in W^{1, n}\left(\Omega, R^{k}\right)$ and $B \subset \Omega$ is a ball. Then

$$
\begin{equation*}
\operatorname{osc}_{B} u \leq 4 \max \left\{\alpha_{1}, \alpha_{2}\right\}, \tag{86}
\end{equation*}
$$

where $\alpha_{1}=\operatorname{osc}_{\partial B} u, \alpha_{2}=\sup _{y \in G \cap B} \inf _{x \in \partial B}|u(x)-u(y)|$.

Proof: The proof is similar to that in [16]. Denote $\alpha=\max \left\{\alpha_{1}, \alpha_{2}\right\}$. Take a point $x_{1} \in \partial B$. Define $z=u-u\left(x_{1}\right)$ and $v=\max \{|z|-2 \alpha, 0\}$. Then $v \in W_{0}^{1, n}(\Omega)$. From the definitions of $\alpha_{1}, \alpha_{2}$, one sees that $v=0$ on $G \cap B$. It follows that $\nabla v(y)=0$ if $y \in G \cap B$. On the complement of $G \cap B, \nabla v=0$ almost everywhere. Therefore $\nabla v=0$ almost everywhere on $B$, and so $v$ must be a constant, which is zero.

We also need the Courant-Lebesgue Lemma.
Lemma 20 Suppose $u \in W^{1, n}\left(\Omega, R^{k}\right)$ and $B(x, r) \subset \Omega, 0<r<1$. Then there is a constant $C>0$ and some $\delta \in\left[\frac{r}{2}, r\right]$ such that

$$
\begin{equation*}
\underset{\partial B(x, \delta)}{O S C} u \leq C K^{1 / n}, \quad \text { where } K=\int_{B(x, r)}|\nabla u|^{n} . \tag{87}
\end{equation*}
$$

Proof: Recall that for $y \in B(x, r),|\nabla u(y)|^{2} \geq \rho^{-2}\left|\nabla_{\theta} u(y)\right|^{2}$, where $\rho=$ $|y-x|$ and $\theta=\frac{y-x}{\rho} \in S^{n-1}$. It follows

$$
\begin{equation*}
\int_{\frac{r}{2}}^{r} \int_{S^{n-1}} \rho^{-1}\left|\nabla_{\theta} u(y)\right|^{n} d \theta d \rho \leq K \tag{88}
\end{equation*}
$$

By Fubini's theorem, there is a $\delta \in\left[\frac{r}{2}, r\right]$ such that

$$
\begin{equation*}
\int_{\frac{r}{2}}^{r} \int_{S^{n-1}} \delta^{-1}\left|\nabla_{\theta} u(y)\right|^{n} d \theta d \rho=\frac{r}{2 \delta} \int_{S^{n-1}}\left|\nabla_{\theta} u(y)\right|^{n} d \theta \tag{89}
\end{equation*}
$$

Since $\delta \leq r$, (88) (89) imply that $\int_{S^{n-1}}\left|\nabla_{\theta} u(y)\right|^{n} d \theta \leq 2 K$. (87) follows from Sobolev embedding theorem $W^{1, n}\left(S^{n-1}, R^{k}\right) \hookrightarrow C^{1 / n}$.

This lemma gives a control of the oscillation of $u$ on the boundary $\partial B(x, \delta)$. Our following step is to estimate the interior oscillation. We need some propositions.

Proposition 21 Suppose $u \in W^{1, n}\left(\Omega, R^{k}\right)$ is conformal and $B \subset \Omega$ is an open subset. Define $D_{\sigma}=B \cap\left\{x:\left|u(x)-u\left(x_{0}\right)\right|<\sigma\right\}$ for $x_{0} \in B \cap G$ and $\sigma>0$. Then

$$
\begin{equation*}
\underset{\sigma \rightarrow 0}{\limsup } \sigma^{-n} \int_{D_{\sigma}}|\nabla u|^{n} \geq n^{\frac{n}{2}-1} \omega_{n-1}, \tag{90}
\end{equation*}
$$

where $\omega_{n-1}$ is the area of the sphere $S^{n-1}$.

Remark 22 This Proposition was proved by Grüter [16] when $n=2$.
Remark 23 Without the assumption that $u$ is conformal, then (90) still holds, with the right hand being replaced by $n^{-1} \omega_{n-1}$.

Proof: We may assume $u\left(x_{0}\right)=0$. For $\epsilon, \sigma>0$, define

$$
\begin{aligned}
& T_{\epsilon}=B \backslash\left\{x:\left|u(x)-\nabla u\left(x_{0}\right)\left(x-x_{0}\right)\right| \leq \frac{\epsilon}{\sqrt{n}}\left|x-x_{0}\right|\right\}, \\
& B_{\epsilon}=B \cap\left\{x:\left|x-x_{0}\right|<r_{\epsilon}\right\} \text {, where } r_{\epsilon}=\frac{\sigma \sqrt{n}}{\left|\nabla u\left(x_{0}\right)\right|+\epsilon} .
\end{aligned}
$$

We claim

$$
B_{\epsilon} \backslash T_{\epsilon} \subset D_{\sigma} \backslash T_{\epsilon} .
$$

Indeed, if $x \in B_{\epsilon} \backslash T_{\epsilon}$, then

$$
\left|u(x)-\nabla u\left(x_{0}\right)\left(x-x_{0}\right)\right| \leq \frac{\epsilon}{\sqrt{n}}\left|x-x_{0}\right| ;
$$

while the conformality condition (8) implies that

$$
\left|\nabla u\left(x_{0}\right)\left(x-x_{0}\right)\right|^{2} \leq \frac{1}{n}\left|\nabla u\left(x_{0}\right)\right|^{2}\left|x-x_{0}\right|^{2} .
$$

Therefore,

$$
|u(x)| \leq \frac{1}{\sqrt{n}}\left(\left|\nabla u\left(x_{0}\right)\right|+\epsilon\right)\left|x-x_{0}\right|<\sigma,
$$

and so $x \in D_{\sigma} \backslash T_{\epsilon}$.
Note that for any $a, b \geq 0, \epsilon>0$, and $p>1$, there holds

$$
a^{p} \geq(1-\epsilon) b^{\epsilon}-\left(\epsilon^{-1}-1\right)\left|a^{p}-b^{p}\right| .
$$

(For a proof, note that $\epsilon b^{p}+\epsilon^{-1}\left|a^{p}-b^{p}\right| \geq 2 b^{p / 2}\left|a^{p}-b^{p}\right|^{1 / 2} \geq\left|a^{p}-b^{p}\right|-$ $\left(a^{p}-b^{p}\right)$.) Using this inequality, we obtain

$$
\begin{align*}
\sigma^{-n} \int_{D_{\sigma}}|\nabla u(x)|^{n} \geq & \sigma^{-n} \int_{B_{\epsilon} \backslash T_{\epsilon}}|\nabla u(x)|^{n} \\
\geq & \sigma^{-n}(1-\epsilon) \int_{B_{\epsilon} \backslash T_{\epsilon}}\left|\nabla u\left(x_{0}\right)\right|^{n}-  \tag{91}\\
& -\sigma^{-n}\left(\epsilon^{-1}-1\right) \int_{B_{\epsilon} \backslash T_{\epsilon}}\left(|\nabla u(x)|^{n}-\left|\nabla u\left(x_{0}\right)\right|^{n}\right) .
\end{align*}
$$

We look at each term in (91) as $\sigma \rightarrow 0$. For the first term,

$$
\begin{align*}
& \sigma^{-n}(1-\epsilon) \int_{B_{\epsilon} \backslash T_{\epsilon}}\left|\nabla u\left(x_{0}\right)\right|^{n}=\sigma^{-n}(1-\epsilon)\left|\nabla u\left(x_{0}\right)\right|^{n} L^{n}\left(B_{\epsilon} \backslash T_{\epsilon}\right)  \tag{92}\\
\rightarrow & \frac{\omega_{n-1}}{n} \sqrt{n^{n}}(1-\epsilon)\left(\frac{\left|\nabla u\left(x_{0}\right)\right|}{\left|\nabla u\left(x_{0}\right)\right|+\epsilon}\right)^{n}\left(1-\Phi^{n}\left(L\left\lfloor T_{\epsilon}, x_{0}\right)\right)\right. \\
= & \frac{\omega_{n-1}}{n} \sqrt{n^{n}}(1-\epsilon)\left(\frac{\left|\nabla u\left(x_{0}\right)\right|}{\left|\nabla u\left(x_{0}\right)\right|+\epsilon}\right)^{n} .
\end{align*}
$$

For the second term in (91), using that $x_{0}$ is a Lebesgue point of $|\nabla u|^{n}$, we have, for fixed $\epsilon$,

$$
\begin{align*}
& \left.\left.\sigma^{-n}\left(\epsilon^{-1}-1\right) \int_{B_{\epsilon} \backslash T_{\epsilon}}| | \nabla u(x)\right|^{n}-\left|\nabla u\left(x_{0}\right)\right|^{n}\right) \mid \\
& \left.\leq\left.\frac{\left(\epsilon^{-1}-1\right) \sqrt{n^{n}}}{\left(\left|\nabla u\left(x_{0}\right)\right|+\epsilon\right)^{n} L^{n}\left(B_{\epsilon}\right)} \int_{B_{\epsilon} \backslash T_{\epsilon}}| | \nabla u(x)\right|^{n}-\left|\nabla u\left(x_{0}\right)\right|^{n} \right\rvert\,  \tag{93}\\
& \rightarrow 0 \text { as } \sigma \rightarrow 0 .
\end{align*}
$$

Now taking the limit in (91) and using (92) (93), we obtain (90).
Theorem 24 Suppose $B \subset \Omega$ is a ball, $x_{0} \in B \cap G$, and $\Sigma>0$ is a number such that

$$
\begin{gather*}
2 n \Lambda \Sigma \leq 1,  \tag{94}\\
\operatorname{dist}\left(u(\partial B), u\left(x_{0}\right)\right)>\Sigma, \tag{95}
\end{gather*}
$$

where $\Lambda$ is as in (2). Then

$$
\begin{equation*}
\int_{B}|\nabla u|^{n} \geq \frac{1}{2} \omega_{n-1} n^{\frac{n}{2}-1} \Sigma^{n} . \tag{96}
\end{equation*}
$$

Corollary 25 Suppose $B \subset \Omega$ is a ball such that

$$
\int_{B}|\nabla u|^{n} \leq \frac{\omega_{n-1}}{2^{\bar{n}+2} n^{n / 2+1} \Lambda^{n}} .
$$

Then for any $x_{0} \in B \cap G$,

$$
\operatorname{dist}\left(u(\partial B), u\left(x_{0}\right)\right) \leq \frac{1}{\omega_{n-1} n^{\frac{n}{2}-1}}\left(\int_{B}|\nabla u|^{n}\right)^{1 / n} .
$$

Proof of Corollary 25: Let

$$
\Sigma=\left(\frac{4}{\omega_{n-1} n^{n / 2-1}} \int_{B}|\nabla u|^{n}\right)^{1 / n}
$$

Then the condition (94) is satisfied, but the conclusion (96) does not hold, therefore, (95) must not hold.
Proof of Theorem 24: For $\sigma \in(0, \Sigma]$, denote

$$
D_{\sigma}=B \cap\left\{x:\left|u(x)-u\left(x_{0}\right)\right|<\sigma\right\} .
$$

Let $\lambda \in C_{0}^{1}(R,[0,1])$ be a function such that $\lambda(t)=0$ for $t \leq 0$. For $\rho \in(0, \Sigma)$, define

$$
\eta=\lambda(\rho-|u|) u .
$$

From (94), we have $\eta \in W_{0}^{1, n}\left(B, R^{k}\right) \cap L^{\infty}$. Multiplying $\eta$ to the equation (1) and integrating by parts, we obtain

$$
\begin{align*}
\int_{D_{\rho}}|\nabla u|^{n} \lambda(\rho & -|u|)-\int_{D_{\rho}}|\nabla u|^{n-2} \lambda^{\prime}(\rho-|u|)\left|\nabla u \cdot \frac{u}{|u|}\right|^{2} \\
& =\int_{D_{\rho}} f(x, u, \nabla u) u \lambda(\rho-|u|) \tag{97}
\end{align*}
$$

Define

$$
\Phi(\rho)=\frac{1}{n} \int_{D_{\rho}}|\nabla u|^{n} \lambda(\rho-|u|) .
$$

Then we have

$$
\begin{equation*}
\Phi^{\prime}(\rho) \geq \frac{1}{n} \int_{D_{\rho}}|\nabla u|^{n} \lambda^{\prime}(\rho-|u|) \tag{98}
\end{equation*}
$$

From the conformality of $u$, it follows that $|\nabla u \cdot u|^{2} \leq \frac{1}{n}|\nabla u|^{2}|u|^{2}$. The property of $\lambda$ implies that

$$
\lambda^{\prime}(\rho-|u|)|u| \leq \rho \lambda^{\prime}(\rho-|u|) .
$$

Therefore, we have that

$$
\begin{gather*}
\int_{D_{\rho}}|\nabla u|^{n-2} \lambda^{\prime}(\rho-|u|)\left|\nabla u \cdot \frac{u}{|u|}\right|^{2} \leq \frac{1}{n} \int_{D_{\rho}}|\nabla u|^{n} \lambda^{\prime}(\rho-|u|)|u|  \tag{99}\\
\leq \frac{1}{n} \rho \int_{D_{\rho}}|\nabla u|^{n} \lambda^{\prime}(\rho-|u|) \leq \rho \Phi^{\prime}(\rho) .
\end{gather*}
$$

Also, we have

$$
\begin{align*}
\int_{D_{\rho}} f(x, u, \nabla u) u \lambda(\rho-|u|) & \leq \Lambda \int_{D_{\rho}}|\nabla u|^{n} \lambda(\rho-|u|)|u| \\
& \leq \Lambda \int_{D_{\rho}}|\nabla u|^{n} \int_{0}^{\rho} \lambda^{\prime}(\sigma-|u|)|u| d \sigma \\
& \leq \Lambda \int_{0}^{\rho} \sigma\left(\int_{D_{\rho}}|\nabla u|^{n} \lambda^{\prime}(\sigma-|u|)|u|\right) d \sigma \\
& \leq n \Lambda \int_{0}^{\rho} \sigma \Phi^{\prime}(\sigma) d \sigma . \tag{100}
\end{align*}
$$

Thus (97) together with (99) and (100) yields

$$
\begin{equation*}
n \Phi(\rho)-\rho \Phi^{\prime}(\rho) \leq n \Lambda \int_{0}^{\rho} \sigma \Phi^{\prime}(\sigma) d \sigma \tag{101}
\end{equation*}
$$

This can be rewritten as

$$
\begin{equation*}
-\left(\frac{\Phi(\rho)}{\rho^{n}}\right)^{\prime} \leq \frac{n \Lambda}{\rho^{n+1}} \int_{0}^{\rho} \sigma \Phi^{\prime}(\sigma) d \sigma \leq n \Lambda \frac{\Phi(\rho)}{\rho^{n}} . \tag{102}
\end{equation*}
$$

This differential inequality implies that $e^{n \Lambda \rho} \frac{\Phi(\rho)}{\rho^{n}}$ is increasing in $\rho$; in particular, $\frac{\Phi(\rho)}{\rho^{n}}$ has a limit as $\rho \rightarrow 0$. Furthermore, for $0<\rho_{1} \leq \rho_{2} \leq \Sigma$, by integrating 102 from $\rho_{1}$ to $\rho_{2}$, we have

$$
\begin{equation*}
\frac{\Phi\left(\rho_{1}\right)}{\rho_{1}^{n}} \leq \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n}}+n \Lambda \int_{0}^{\rho_{2}} \frac{\Phi(\rho)}{\rho^{n}} d \rho \tag{103}
\end{equation*}
$$

The second term of (103) can be estimated by integration by parts and using (101)

$$
\begin{align*}
\int_{0}^{\rho_{2}} \frac{\Phi(\rho)}{\rho^{n}} d \rho & \leq \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n-1}}+\int_{0}^{\rho_{2}} \rho\left(-\frac{\Phi(\rho)}{\rho^{n}}\right)^{\prime} d \rho \\
& \leq \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n-1}}+n \Lambda \int_{0}^{\rho_{2}} \frac{1}{\rho^{n}} \int_{0}^{\rho} \sigma \Phi^{\prime}(\sigma) d \sigma d \rho  \tag{104}\\
& \leq \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n-1}}+n \Lambda \int_{0}^{\rho_{2}} \frac{\Phi(\rho)}{\rho^{n-1}} d \rho \\
& \leq \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n-1}}+n \Lambda \rho_{2} \int_{0}^{\rho_{2}} \frac{\Phi(\rho)}{\rho^{n}} d \rho
\end{align*}
$$

From the assumption, $n \Lambda \rho_{2} \leq n \Lambda \Sigma \leq \frac{1}{2}$. Thus it follows that from (104)

$$
\begin{equation*}
\int_{0}^{\rho_{2}} \frac{\Phi(\rho)}{\rho^{n}} d \rho \leq 2 \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n-1}} \tag{105}
\end{equation*}
$$

Now (103) and (105) imply that

$$
\begin{equation*}
\frac{\Phi\left(\rho_{1}\right)}{\rho_{1}^{n}} \leq \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n}}+2 n \Lambda \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n-1}} \leq 2 \frac{\Phi\left(\rho_{2}\right)}{\rho_{2}^{n}} . \tag{106}
\end{equation*}
$$

Given $\epsilon>0$, we choose $\lambda(t)$ with the additional property that $\lambda(t)=1$ for $t \geq \epsilon \rho_{1}$. Then it easy to see

$$
\begin{equation*}
\frac{1}{n} \int_{D_{\rho_{1}(1-\epsilon)}}|\nabla u|^{n} \leq \Phi\left(\rho_{1}\right), \quad \Phi\left(\rho_{2}\right) \leq \frac{1}{n} \int_{D_{\rho_{2}}}|\nabla u|^{n} . \tag{107}
\end{equation*}
$$

Apply (106) with $\rho_{2}=\Sigma$, and use (107), we then obtain

$$
\begin{equation*}
\frac{\int_{D_{\rho_{1}(1-\epsilon)}}|\nabla u|^{n}}{\rho_{1}^{n}} \leq 2 \frac{\int_{D_{\Sigma}}|\nabla u|^{n}}{\Sigma^{n}} \leq 2 \frac{\int_{B}|\nabla u|^{n}}{\Sigma^{n}} . \tag{108}
\end{equation*}
$$

Let $\rho_{1} \rightarrow 0$ in (108) and apply Proposition 21, and then send $\epsilon \rightarrow 0$. (96) then follows.

Proof of Theorem 17. We divide the proof into several steps. The first step, showing the continuity of $u$, is the essential one. The other steps are standard.
(a). $u$ is continuous. Fix $x_{0} \in \Omega$. Choose $R>0$ small enough such that $B\left(x_{0}, R\right) \subset \Omega$ and $\int_{B\left(x_{0}, R\right)}|\nabla u|^{n} \leq \frac{\omega_{n-1}}{2^{\bar{n}+2} n^{n / 2+1} \Lambda^{n}}$. By the Courant-Lebesgue Lemma 20, there is a $\delta \in\left[\frac{R}{2}, R\right]$ such that

$$
\underset{\partial B\left(x_{0}, \delta\right)}{\mathrm{OSC}} u \leq C\left(\int_{B\left(x_{0}, R\right)}|\nabla u|^{n}\right)^{1 / n} \equiv \alpha_{1}(R) .
$$

Because $\int_{B\left(x_{0}, \delta\right)}|\nabla u|^{n} \leq \int_{B\left(x_{0}, R\right)}|\nabla u|^{n}$, by Corollary 25, for any $x_{0} \in B \cap$ $G, B=B\left(x_{0}, \delta\right)$,

$$
\text { dist }\left(u(\partial B), u\left(x_{0}\right)\right) \leq \frac{1}{\omega_{n-1} n^{\frac{n}{2}-1}}\left(\int_{B\left(x_{0}, R\right)}|\nabla u|^{n}\right)^{1 / n} \equiv \alpha_{2}(R) .
$$

Now Lemma 19 implies that

$$
\begin{equation*}
\operatorname{osc}_{B\left(x_{0}, \delta\right)} u \leq 4 \max \left\{\alpha_{1}(R), \alpha_{2}(R)\right\} \rightarrow 0 \text { as } R \rightarrow 0 . \tag{109}
\end{equation*}
$$

In particular, $\operatorname{osc}_{B\left(x_{0}, R / 2\right)} u \rightarrow 0$ as $R \rightarrow 0$. So $u$ is continuous at $x_{0}$.
(b). $u$ is $C^{\beta}$ for some $\beta \in(0,1)$. We claim that if $x_{0} \in W$ and $R>0$ such that $B\left(x_{0}, R\right) \subset \Omega$ and $\operatorname{osc}_{B\left(x_{0}, R\right)} u \Lambda<1$, then there is a number $\tau \in(0,1)$ so that for every ball $B(x, r) \subset B\left(x_{0}, R\right)$

$$
\begin{equation*}
\int_{B(x, r / 2)}|\nabla u|^{n} \leq \tau \int_{B(x, r)}|\nabla u|^{n} . \tag{110}
\end{equation*}
$$

By iteration, we have for some constants $C$ and $\gamma \in(0,1)$,

$$
\int_{B(x, r)}|\nabla u|^{n} \leq C r^{\gamma} \text { for } x \in B\left(x_{0}, R / 2\right) \text { and } r \in(0, R / 2) .
$$

That $u$ is $C^{\beta}$ for some $\beta \in(0,1)$ follows from Morrey's Lemma [23][3.5.2]. The proof of (110) is a standard "hole-filling" method. Let $\bar{u}$ be the mean value of $u$ on $B(x, r) \backslash B(x, r / 2)$ and $\eta \in C_{0}^{1}(B(x, r),[0,1])$ be a cut-off function such that $\eta=1$ on $B(x, r / 2)$ and $|\nabla \eta| \leq 3 / r$. Take $\phi=(u-\bar{u}) \eta$, then $\phi \in W_{0}^{1, n}\left(B(x, r), R^{k}\right)$. Multiply $\phi$ to the equation (1) and integrate. We then get

$$
\begin{align*}
& \left.\left|\int_{B(x, r)} \eta\right| \nabla u\right|^{n}+\int_{B(x, r)}|\nabla u|^{n-2} \nabla u \nabla \eta(u-\bar{u}) \mid  \tag{111}\\
= & \left|\int_{B(x, r)} f(x, u, \nabla u) \eta(u-\bar{u})\right| \\
\leq & \operatorname{osc}_{B\left(x_{0}, R\right)} u \Lambda \int_{B(x, r)}|\nabla u|^{n} .
\end{align*}
$$

By Hölder and Poincare's inequalities, we estimate the second term of (111),

$$
\begin{align*}
& \left.\left|\int_{B(x, r)}\right| \nabla u\right|^{n-2} \nabla u \nabla \eta(u-\bar{u}) \mid \\
\leq & C_{1}\left(\int_{B(x, r) \backslash B(x, r / 2)}|\nabla u|^{n}\right)^{(n-1) / n}\left(\frac{1}{r} \int_{B(x, r) \backslash B(x, r / 2)}|u-\bar{u}|^{n}\right)^{1 / n} \\
\leq & C_{2} \int_{B(x, r) \backslash B(x, r / 2)}|\nabla u|^{n}, \tag{112}
\end{align*}
$$

where $C_{1}$ and $C_{2}$ depend only on $n$ and $k$. Put (112) back to (111) and note the property of $\eta$. It follows

$$
\int_{B(x, r / 2)}|\nabla u|^{n}-C_{2} \int_{B(x, r) \backslash B(x, r / 2)}|\nabla u|^{n} \leq \operatorname{osc}_{B\left(x_{0}, R\right)} u \Lambda \int_{B(x, r)}|\nabla u|^{n}
$$

or

$$
\left(C_{2}+1\right) \int_{B(x, r / 2)}|\nabla u|^{n} \leq\left(C_{2}+\operatorname{osc}_{B\left(x_{0}, R\right)} u \Lambda\right) \int_{B(x, r)}|\nabla u|^{n}
$$

Let $\tau=\left(C_{2}+\right.$ osc $\left._{B\left(x_{0}, R\right)} u \Lambda\right) /\left(C_{2}+1\right)$. Then (43) follows.
(c) . $u$ is $C^{1, \alpha}$ for some $\alpha \in(0,1)$. For the proof of $C^{1, \alpha}$ regularity based on $C^{\beta}$, we refer [17] or [13].
(d). $u$ is continuous up to $\partial \Omega$. This was proved in [24] [Theorem 4.1].

Acknowledgement. The authors wish to thank M. Fuchs for bring up their attention to his results in [10][14][9] and other works on surfaces with prescribed mean curvature vectors and on n-harmonic maps.

## References

[1] E. Acerbi and N. Fusco, Semicontinuity problems in the calculus of variations. Arch. Rat. Mech. Anal. 86 (1984) 125-145.
[2] F. Almgren, Optimal isoperimetric inequalities. Indiana University Mathematical Journal, No. 3, Vol. 35 (1986) 451-547.
[3] A. Ambrosetti and P. Rabinowitz, Dual variational methods in critical point theory and applications. J. Funct. anal. 14 (1973) 349-381.
[4] F. Bethuel, On the singular set of stationary harmonic maps. Man. Math. 78 (1993), no.4, 417-443.
[5] F. Bethuel, Un résultat de régularité pour les solutions de l'équation des surfaces á courbure moyenne prescrite. C. R. Acad. Sci. Paris, t. 314, Série I (1992) 1003-1007.
[6] H. Brezis and J.-M. Coron, Multiple solutions of $H$-systems and Rellich's conjecture. Comm. Pure Appl. Math. 37 (1984) 149-187.
[7] H. Brezis and L. Nirenberg, Positive solutions of nonlinear elliptic equations involving critical Sobolev exponents. Comm. Pure \& Appl. Math. Vol. XXXVI (1983) 437-477.
[8] B. Dacorogna, Direct methods in the calculus of variations. Springer, Berlin-Heildelberg-New York (1989).
[9] F. Duzaar and M. Fuchs, On removable singularities of p-harmonic maps, Ann. Inst. Henri Poincare, Vol. 7, 5 (1990) 385-405.
[10] F. Duzaar and M. Fuchs, Existenz und Regularität von Hyperflächen mit vorgeschriebener mittlerer krummung. Analysis no.2-3, 10 (1990) 193-230.
[11] L. C. Evans, Partial regularity for stationary harmonic maps into spheres. Arch. Rat. Mech. Anal. 116 (1991) 101-113.
[12] H. Federer, Geometric Measure Theory. Grundelhren 153, Berlin-Heidelberg-New York, 1969.
[13] M. Fuchs, p-harmonic maps obstacle problems, Part 1: Partial regularity Theory. Annali Mat. Pura Appl. 4, 156(1990) 127-158.
[14] M. Fuchs, The blow-up of p-harmonic maps. Manuscripta Math. 81 (1993) 89-94.
[15] M. Giaquinta, Multiple integrals in the calculus of variations and nonlinear elliptic systems. Ann. Math. Studies 105, Princeton Univ. Press, Princeton (1983).
[16] M. Grüter, Regularity of weak H-surfaces. J. für reine u. angew. Math. 329 (1981) 1-15.
[17] R. Hardt and F.-H. Lin, Mappings minimizing the $L^{p}$ norm of gradient. Comm. Pure Appl. Math. Vol. XL (1987) 555-588.
[18] R. Hardt, F.-H. Lin and L. Mou, Strong convergence of p-harmonic mappings. Progress in partial differential equations: the Metz surveys 3, edited by M. Chipot, J.S.J. Paulin and I. Shafris. Longman Scientific Technical, 1994.
[19] F. Hélein, Regularite des applications faiblement harmoniques entre une surface et variete riemannienne, CRAS, Paris 312 (1991) 591-596.
[20] E. Heinz, Ein Regularitätssatz für schwache Lösungen nichtlinearer elliptischer Systeme. Nachr. Akad. Wiss. Göttingen, Math. Phys. Kl. II, 1 (1977).
[21] S. Hildebrandt, On the Plateau problem for surfaces of constant mean curvature. Comm. Pure Appl. Math. 23 (1970) 97-114.
[22] J. Lewis, Smoothness of certain degenerate elliptic systems. Proc. Amer. Math. Soc. 80 (1980) 259-265.
[23] C. B. Morrey, Multiple integrals in the calculus of variations. Berlin-Heidelberg-New York, 1966.
[24] L. Mou and P. Yang, Regularity of n-harmonic maps. Journal of Geometric Analysis. 6(1996), no. 1, 91-112.
[25] R. Schoen, Analytic aspects of the harmonic map problems, in "Seminar on Nonlinear PDE" edited by S. S. Chern, Springer-Verlag, New York, Berlin, Heidelberg, 1984.
[26] R. Schoen and K. Uhlenbeck, A regularity theory for harmonic maps. Jour. Diff. Geom. 17 (1982) 307-335.
[27] Steffen, On the existence of surfaces with prescribed mean curvature and boundary. Math. Z. 146 (1976) 113-135.
[28] M. Struwe, Plateau's problem and the calculus of variations. Princeton University Press, Mathematical Notes 35 (1988).
[29] M. Struwe, Non-uniqueness in the Plateau problem for surfaces of constant mean curvature. Arch. Rat. Mech. Anal. 93 (1986) 135-157.
[30] T. Toro and C.Y. Wang, Compactness properties of weakly p-harmonic mapping into homogeneous spaces. Ind. Univ. Math. Jour. Vol. 44, No. 1 (1995).
[31] H. C. Wente, An existence theorem for surface of constant mean curvature. J. Math. Anal. Appl. 26 (1969), 318-344.
[32] H. C. Wente, The Dirichlet problem with a volume constraint. Man. Math. 11 (1974) 141-157.
[33] H. C. Wente, Large solutions to volume constrained Plateau problem. Arch. Rat. Mech. Anal. 75 (1980) 59-77.
L.M.: Department of Mathematics, University of Iowa, Iowa City, IA 52242 mou@math.uiowa.edu Current address: Department of Mathematics, Bradley University, Peoria, IL 61625 mou@bradley.bradley.edu
P.Y.: Department of Mathematics, University of Southern California, Los Angeles, CA 90089
pyang@math.usc.edu

